

Lecture 1

- Go over syllabus

Section 4.4.3

Last semester, finished with discussion of spin.
Pick up from there.

Recall electrons have $S = \frac{1}{2}$

Two spin states $|m = +\frac{1}{2}\rangle \equiv |\uparrow\rangle$
 $|m = -\frac{1}{2}\rangle \equiv |\downarrow\rangle$

General state $|z\rangle = \cos\theta |\uparrow\rangle + e^{i\phi} \sin\theta |\downarrow\rangle$

Convenient to use raising and lowering operators.

Generally $J_{\pm} |j, m\rangle = \hbar \sqrt{(j \mp m)(j \pm m + 1)} |j, m \pm 1\rangle$

So for $J = S = \frac{1}{2}$, have

Also:

$$S_z |\uparrow\rangle = \frac{\hbar}{2} |\uparrow\rangle \quad S_+ |\uparrow\rangle = 0 \quad S_+ |\downarrow\rangle = \hbar |\uparrow\rangle$$

$$S_z |\downarrow\rangle = -\frac{\hbar}{2} |\downarrow\rangle \quad S_- |\uparrow\rangle = \hbar |\downarrow\rangle \quad S_- |\downarrow\rangle = 0$$

What we're going to be talking about is
how QM works with multiple particles

We'll start with spin. Let's say we have
two spin- $\frac{1}{2}$ particles... say electron and proton
of a hydrogen atom.

How do we write state of entire atom?

Now four states: $\{ | \uparrow \uparrow \rangle, | \uparrow \downarrow \rangle, | \downarrow \uparrow \rangle, | \downarrow \downarrow \rangle \}$ make basis

where $| \uparrow \downarrow \rangle =$ electron \uparrow , proton \downarrow

But we'd like to know about total angular momentum

$$\vec{S} = \vec{S}^{(1)} + \vec{S}^{(2)}$$

adding operators is nothing special, just apply each one to where it belongs.

For instance, if $| \psi \rangle = | m_1, m_2 \rangle = | m_1 \rangle | m_2 \rangle$

$$\begin{aligned} \text{then } S_z | \psi \rangle &= (S_z^{(1)} + S_z^{(2)}) | m_1 \rangle | m_2 \rangle \\ &= | m_2 \rangle S_z^{(1)} | m_1 \rangle + | m_1 \rangle S_z^{(2)} | m_2 \rangle \\ &= \hbar(m_1 + m_2) | m_1, m_2 \rangle \end{aligned}$$

Since our basis states have definite values of m_1 & m_2 , they have definite values of S_z :

$$\begin{aligned} | \uparrow \uparrow \rangle &\rightarrow m = 1 \\ | \uparrow \downarrow \rangle &\rightarrow m = 0 \\ | \downarrow \uparrow \rangle &\rightarrow m = 0 \\ | \downarrow \downarrow \rangle &\rightarrow m = -1 \end{aligned}$$

What does this say about total spin of atom?

We know $m=1$ shows up for $L=1$ systems

Suggest total spin $S=1$, make sense since $\frac{1}{2} + \frac{1}{2} = 1$

But, $S=1$ system has three states: $m=-1, 0, 1$

We have an extra $m=0$. What's that about?

We can check for sure. Start with $|11\rangle$ state,

Is this an eigenstate of S^2 ?

$$\text{Use } S^2 = (\vec{S}_1 + \vec{S}_2)^2 = S_1^2 + S_2^2 + 2\vec{S}_1 \cdot \vec{S}_2$$

$$\text{and } \vec{S}_1 \cdot \vec{S}_2 = S_x^{(1)} S_x^{(2)} + S_y^{(1)} S_y^{(2)} + S_z^{(1)} S_z^{(2)}$$

$$\text{Remember } S_x = \frac{1}{2}(S_+ + S_-)$$

$$S_y = \frac{1}{2i}(S_+ - S_-)$$

$$\text{So } S_x^{(1)} S_x^{(2)} = \frac{1}{4} (S_+^{(1)} S_+^{(2)} + S_-^{(1)} S_-^{(1)} + S_+^{(1)} S_-^{(2)} + S_-^{(1)} S_+^{(1)})$$

$$S_y^{(1)} S_y^{(2)} = -\frac{1}{4} (S_+^{(1)} S_+^{(2)} + S_-^{(1)} S_-^{(1)} - S_+^{(1)} S_-^{(1)} - S_-^{(1)} S_+^{(1)})$$

$$\text{and } \vec{S}_1 \cdot \vec{S}_2 = \frac{1}{2} (S_+^{(1)} S_-^{(1)} + S_-^{(1)} S_+^{(2)}) + S_z^{(1)} S_z^{(2)}$$

$$\text{So, apply } S^2 \text{ to } |11\rangle: \quad S_1^2 |11\rangle = (S_1^2 |1\rangle) |1\rangle = \hbar^2 (s_1)(s_1+1) |11\rangle \\ = \hbar^2 \frac{1}{2} \cdot \frac{3}{2} |11\rangle \\ S_2^2 |11\rangle = \frac{3}{4} \hbar^2 |11\rangle = \frac{3}{4} \hbar^2 |11\rangle$$

$$2\vec{S}_1 \cdot \vec{S}_2 |11\rangle = 0 + 2S_z^{(1)} S_z^{(2)} |11\rangle = 2\hbar^2 \cdot \frac{1}{2} \cdot \frac{1}{2} |11\rangle = \frac{1}{2} \hbar^2 |11\rangle$$

When $S_+^{(1)}$ hits $|1\rangle$, get zero

$$\begin{aligned} \text{Thus } S^2 | \uparrow \uparrow \rangle &= \hbar^2 \left(\frac{3}{4} + \frac{3}{4} + \frac{1}{2} \right) | \uparrow \uparrow \rangle \\ &= 2\hbar^2 | \uparrow \uparrow \rangle = \hbar^2 (1)(1+1) | \uparrow \uparrow \rangle \end{aligned}$$

Thus $| \uparrow \uparrow \rangle$ has total spin = 1

and $m=1$, so write $| \uparrow \uparrow \rangle = | 1 1 \rangle$

To sort out what's going on with $m=0$ states,
just apply S_- to $| \uparrow \uparrow \rangle$

$$\text{Know } S_- | 1 1 \rangle = \hbar \sqrt{2} | 1 0 \rangle$$

$$\begin{aligned} \text{But also } S_- | \uparrow \uparrow \rangle &= (S_-^{(1)} + S_-^{(2)}) | \uparrow \uparrow \rangle \\ &= (S_-^{(1)} | \uparrow \rangle) | \uparrow \rangle + | \uparrow \rangle (S_-^{(2)} | \uparrow \rangle) \\ &= \hbar (| \downarrow \uparrow \rangle + | \uparrow \downarrow \rangle) \end{aligned}$$

$$\text{indicating that } | 1 0 \rangle = \frac{1}{\sqrt{2}} (| \uparrow \downarrow \rangle + | \downarrow \uparrow \rangle)$$

symmetric combination of our basis states.

$$\text{Easy to see that } | \downarrow \downarrow \rangle = | 1 -1 \rangle$$

So, we do get the three $S=1$ states,

But we have four states total. Remaining orthogonal
one is $| 0 0 \rangle = \frac{1}{\sqrt{2}} (| \uparrow \downarrow \rangle - | \downarrow \uparrow \rangle)$

Evidently, this is not an $S=1$ state.
See what it is:

$$\text{Apply } S^2 = S_1^2 + S_2^2 + 2S_{1z}^{(1)} S_{2z}^{(2)} + S_+^{(1)} S_-^{(2)} + S_-^{(1)} S_+^{(2)}$$

$$S_1^2 |4\rangle = S_2^2 |4\rangle = \frac{3}{4} \hbar^2 |4\rangle$$

$$\begin{aligned} S_{1z}^{(1)} S_{2z}^{(2)} |4\rangle &= \frac{1}{\sqrt{2}} (S_{1z}^{(1)} |\uparrow\rangle S_{2z}^{(2)} |\downarrow\rangle - S_{1z}^{(1)} |\downarrow\rangle S_{2z}^{(2)} |\uparrow\rangle) \\ &= \frac{1}{\sqrt{2}} \left(-\frac{\hbar^2}{4} |\uparrow\downarrow\rangle - \left(-\frac{\hbar^2}{4}\right) |\downarrow\uparrow\rangle \right) \\ &= -\frac{\hbar^2}{4} |4\rangle \end{aligned}$$

$$\begin{aligned} S_+^{(1)} S_-^{(2)} |4\rangle &= \frac{1}{\sqrt{2}} (S_+^{(1)} |\uparrow\rangle S_-^{(2)} |\downarrow\rangle - S_+^{(1)} |\downarrow\rangle S_-^{(2)} |\uparrow\rangle) \\ &= \frac{1}{\sqrt{2}} (0 - \hbar^2 |\uparrow\downarrow\rangle) = -\frac{\hbar^2}{\sqrt{2}} |\uparrow\downarrow\rangle \end{aligned}$$

$$\begin{aligned} S_-^{(1)} S_+^{(2)} |4\rangle &= \frac{1}{\sqrt{2}} (S_-^{(1)} |\uparrow\rangle S_+^{(2)} |\downarrow\rangle - S_-^{(1)} |\downarrow\rangle S_+^{(2)} |\uparrow\rangle) \\ &= \frac{1}{\sqrt{2}} (\hbar^2 |\downarrow\uparrow\rangle - 0) = \frac{\hbar^2}{\sqrt{2}} |\downarrow\uparrow\rangle \end{aligned}$$

$$\begin{aligned} \text{So } S^2 |4\rangle &= \left(\frac{3}{4} \hbar^2 + \frac{3}{4} \hbar^2 - \frac{1}{4} \hbar^2 \right) |4\rangle + \frac{\hbar^2}{\sqrt{2}} (-|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \\ &= \hbar^2 |4\rangle - \hbar^2 |4\rangle \\ &= 0 \end{aligned}$$

Note that $0 = \hbar^2 (0)(0+1)$, so $|4\rangle$ is eigenstate with $S=0$ (and $m=0$)

So in fact, there are four possible spin states for a hydrogen atom. could have $S=1, m=0, \pm 1$ or $S=0, m=0$