

## Notes on the Brydges-Kennedy-Abdesselam-Rivasseau forest interpolation formula

There are many instances in mathematical physics where one tries to understand joint probability measures for a collection of random variables  $X_1, \dots, X_N$ , with  $N$  large, of the form

$$e^{-\sum_{i=1}^N V(x_i)} d\mu_C(x)$$

where  $d\mu_C$  is a Gaussian measure on  $\mathbb{R}^N$ . The dependence between these random variables is entirely due to the Gaussian measure which, in general, is given by covariances

$$C_{ij} = \text{cov}(X_i, X_j)$$

which do not vanish for  $i \neq j$ . A typical procedure one uses in this type of problem is to try to interpolate between the given covariance matrix  $C$  and the covariance obtained by killing the off-diagonal entries. The outcome is what is called a cluster expansion in the constructive field theory literature. The first such expansions appeared in the context of the construction of the  $\phi^4$  model in 2 dimensions and in the infinite volume limit. They are due to Glimm, Jaffe and Spencer [11, 12, 10]. See also [21, 17] for a simpler presentation on the  $\mathbb{Z}^d$  lattice, instead of the continuum model. Later on, a new simpler interpolation/expansion method was introduced for the same purpose by Brydges, Battle and Federbush [7, 4, 5]. The Brydges-Kennedy-Abdesselam-Rivasseau or BKAR formula offers yet another simpler interpolation formula. It allows the resummation of QFT Feynman diagrams according to spanning subtrees. It can also be used in order to prove the Penrose-Rota inequality for Mayer expansion coefficients. What follows is a self-contained presentation of the BKAR identity which is nothing more than a sophisticated form of the Fundamental Theorem of Calculus. The BKAR formula was dubbed the “constructive Swiss knife” in [19].

Let us consider a finite set  $E \neq \emptyset$ , and let us denote by  $E^{(2)}$  the set unordered pairs  $\{a, b\}$ , where  $a$  and  $b$  are any distinct elements in  $E$ . Of course  $|E^{(2)}| = \binom{|E|}{2}$ . We will consider the space  $\mathbb{R}^{E^{(2)}}$  of multiplets  $s = (s_l)_{l \in E^{(2)}}$  indexed by pairs  $l \in E^{(2)}$ , and functions defined on a particular compact convex set  $\mathcal{K}_E$  in this space. Let  $\Pi_E$  denote the set of partitions of  $E$ . For any partition  $\pi = \{X_1, \dots, X_q\}$  in  $\Pi_E$  we associate a vector  $v_\pi = (v_{\pi, l})_{l \in E^{(2)}}$  defined as

$$v_{\pi, l} = \mathbb{1}_{\{\exists i, 1 \leq i \leq q, l \subset X_i\}} \cdot$$

Now  $\mathcal{K}_E$  is by definition the convex hull of the vectors  $v_\pi$ , for  $\pi \in \Pi_E$ . It is easy to see that  $\mathcal{K}_E$  affinely generates  $\mathbb{R}^{E^{(2)}}$ . Indeed, let  $\hat{\circ}$  be the partition entirely made of singletons, and for any pair  $l \in E^{(2)}$  let  $\hat{l}$  denote the partition made of the two element set  $l$  and the singletons  $\{a\}$ , for  $a \in E \setminus l$ . Then, the vectors  $v_{\hat{l}} - v_{\hat{\circ}}$ , for  $l \in E^{(2)}$  form a basis of the vector space  $\mathbb{R}^{E^{(2)}}$ . As a result, the open domain  $\Omega_E = \mathring{\mathcal{K}}_E$  is nonempty, and  $\mathcal{K}_E$  is equal to the closure  $\bar{\Omega}_E$ . Let  $C^k(\bar{\Omega}_E)$  denote the usual space of functions of class  $C^k$  on the domain  $\Omega_E$

which, together with their derivatives up to order  $k$ , admit uniformly continuous extensions to the closure  $\mathcal{K}_E = \bar{\Omega}_E$  (see, e.g., [3]). It is easy to see that for any  $f$  in  $C^1(\bar{\Omega}_E)$  the Fundamental Theorem of Calculus

$$f(t) = f(s) + \int_0^1 du \frac{d}{du} f((1-u)t + us)$$

holds for any multiplsets  $s$  and  $t$  in  $\mathcal{K}_E$ , even on the boundary. This will be used repeatedly in the following.

Now a simple graph with vertex set  $E$  can be thought of as a subset of the complete graph  $E^{(2)}$ . A forest  $\mathfrak{F}$  is a graph with no circuits, and it is made of a vertex-disjoint collection of trees. Let  $\mathfrak{F}$  be a forest, and let  $\vec{h} = (h_l)_{l \in \mathfrak{F}}$  be a vector of real parameters indexed by the edges  $l$  in the forest  $\mathfrak{F}$ . To such data we canonically associate a multiplset  $s(\mathfrak{F}, \vec{h}) = (s(\mathfrak{F}, \vec{h})_l)_{l \in E^{(2)}}$  in  $\mathbb{R}^{E^{(2)}}$  as follows. Let  $a$ , and  $b$  be two distinct elements in  $E$ . If  $a$  and  $b$  belong to two distinct connected components of the forest  $\mathfrak{F}$ , then  $s(\mathfrak{F}, \vec{h})_{\{a,b\}} = 0$ . Otherwise let, by definition,  $s(\mathfrak{F}, \vec{h})_{\{a,b\}} = \min_l h_l$  where  $l$  belongs to the unique path in the forest  $\mathfrak{F}$  joining  $a$  to  $b$ . We are now ready to state the BKAR formula.

**Theorem 1** [8, 1] *Let  $f \in C^{|E|-1}(\bar{\Omega}_E)$ , and let  $1 \in \mathbb{R}^{E^{(2)}}$  denote the multiplset with all entries equal to one. This is also the same as  $v_{\hat{1}}$  where  $\hat{1}$  is the single block partition  $\{E\}$ . We then have*

$$f(1) = \sum_{\mathfrak{F} \text{ forest}} \int_{[0,1]^{\mathfrak{F}}} d\vec{h} \frac{\partial^{|\mathfrak{F}|} f}{\prod_{l \in \mathfrak{F}} \partial s_l} (s(\mathfrak{F}, \vec{h}))$$

where the sum is over all forests  $\mathfrak{F}$  with vertex set  $E$ , the notation  $d\vec{h}$  is for the Lebesgue measure on the set of parameters  $[0,1]^{\mathfrak{F}}$ , the partial derivatives of  $f$  are with respect to the entries indexed by the pairs belonging to  $\mathfrak{F}$ , and the evaluation of these derivatives is at the  $\vec{h}$  dependent point  $s(\mathfrak{F}, \vec{h})$ . Such points belong to  $\mathcal{K}_E$ .

Note that the empty forest always occurs and its contribution is  $f(0) = f(v_{\emptyset})$ . There are many proofs of this result [8, 1, 9], but we believe the most natural and most easily generalizable is the one given in [2, §2]. It is recalled here for the sake of completeness. We will first prove an ordered forest analog of Theorem 1. A possibly empty sequence  $F = (l_1, \dots, l_p)$  of pairs  $l_i$  in  $E^{(2)}$  is called an ordered forest or o-forest if the corresponding set  $\mathfrak{F} = \{l_1, \dots, l_p\}$  is a forest. Let us denote by  $\Delta_p$  the simplex  $\{\vec{\rho} \in \mathbb{R}^p | 1 \geq \rho_1 \geq \dots \geq \rho_p \geq 0\}$ . Given a vector of parameters  $\vec{\rho} = (\rho_1, \dots, \rho_p)$  in  $\Delta_p$  we also define the multiplset  $t(F, \vec{\rho}) = (t(F, \vec{\rho})_l)_{l \in E^{(2)}}$  as follows. Let  $l = \{a, b\}$ , with  $a$  and  $b$  distinct elements in  $E$ . If  $a$  and  $b$  fall in distinct connected components of the full forest  $\{l_1, \dots, l_p\}$  then we set  $t(F, \vec{\rho})_l = 0$ . Else we let  $t(F, \vec{\rho})_l = \rho_q$  where  $q$  is the smallest index  $1 \leq q \leq p$  such that  $a$  and  $b$  are connected by the subforest  $\{l_1, l_2, \dots, l_q\}$ . The following important property is an easy consequence of this definition.

**Proposition 1** For any  $o$ -forest  $F = (l_1, \dots, l_p)$  and vector  $\vec{\rho} \in \Delta_p$ , the multiplet  $t(F, \vec{\rho})$  belongs to  $\mathcal{K}_E = \bar{\Omega}_E$ .

**Proof:** Indeed one can write the convex combination

$$t(F, \vec{\rho}) = (1 - \rho_1)v_{\hat{0}} + (\rho_1 - \rho_2)v_{\pi_1} + \dots + (\rho_{r-1} - \rho_r)v_{\pi_{r-1}} + \rho_r v_{\pi_r}$$

where  $\pi_q$  is the partition of connected components of the forest  $\{l_1, l_2, \dots, l_q\}$ . Note that we used the fact that  $\pi_0 = \hat{0}$ .  $\blacksquare$

**Theorem 2** Under the same hypotheses as in the previous theorem one has

$$f(1) = \sum_{\substack{F=(l_1, \dots, l_p) \\ o\text{-forest}}} \int_{\Delta_p} d\vec{\rho} \frac{\partial^p f}{\partial s_{l_1} \dots \partial s_{l_p}} (t(F, \vec{\rho}))$$

where the summation allows all possible lengths  $p$  for the  $o$ -forest  $F$ , including  $p = 0$ .

**Proof:** We will prove by induction on  $r \geq 1$  that

$$\begin{aligned} f(1) &= \sum_{p < r} \sum_{\substack{F=(l_1, \dots, l_p) \\ o\text{-forest}}} \int_{\Delta_p} d\vec{\rho} \frac{\partial^p f}{\partial s_{l_1} \dots \partial s_{l_p}} (t(F, \vec{\rho})) \\ &+ \sum_{\substack{F=(l_1, \dots, l_r) \\ o\text{-forest}}} \int_{\Delta_r} d\vec{\rho} \frac{\partial^r f}{\partial s_{l_1} \dots \partial s_{l_r}} (\hat{t}(F, \vec{\rho})) \end{aligned} \quad (1)$$

where  $\hat{t}(F, \vec{\rho})_l$  is defined in the same way as  $t(F, \vec{\rho})_l$  except that if  $l$  does not fall inside a connected component of  $\{l_1, \dots, l_r\}$  one sets  $\hat{t}(F, \vec{\rho})_l$  equal to the last parameter  $\rho_r$  instead of zero. Note that  $\hat{t}(F, \vec{\rho})$  is still in the convex  $\mathcal{K}_E$ , since

$$\begin{aligned} \hat{t}(F, \vec{\rho}) &= t(F, \vec{\rho}) + \rho_r(v_{\hat{1}} - v_{\pi_r}) \\ &= (1 - \rho_1)v_{\hat{0}} + (\rho_1 - \rho_2)v_{\pi_1} + \dots + (\rho_{r-1} - \rho_r)v_{\pi_{r-1}} + \rho_r v_{\hat{1}} \cdot \end{aligned} \quad (2)$$

Now by the Fundamental Theorem of Calculus,

$$\begin{aligned} f(1) &= f(v_{\hat{1}}) = f(v_{\hat{0}}) + \int_0^1 d\rho_1 \frac{d}{d\rho_1} f(\rho_1 v_{\hat{1}} + (1 - \rho_1)v_{\hat{0}}) \\ &= f(0) + \sum_{l_1 \in E^{(2)}} \int_0^1 d\rho_1 \frac{\partial f}{\partial s_{l_1}}(\rho_1 v_{\hat{1}}) \end{aligned}$$

which is Eq. (1) for  $r = 1$ , as can be checked from the given definitions. Suppose (1) has been proven for  $r \geq 1$ , and consider the integrand

$$\frac{\partial^r f}{\partial s_{l_1} \dots \partial s_{l_r}} (\hat{t}(F, \vec{\rho}))$$

of any particular term in the second sum. Using (2) and introducing a new parameter  $\rho_{r+1}$ , we rewrite the argument of the derivative of  $f$  as

$$\begin{aligned} \hat{t}(F, \vec{\rho}) &= (1 - \rho_1)v_{\hat{0}} + (\rho_1 - \rho_2)v_{\pi_1} + \cdots \\ &+ (\rho_{r-1} - \rho_r)v_{\pi_{r-1}} + (\rho_r - \rho_{r+1})v_{\pi_r} + \rho_{r+1}v_{\hat{1}} \Big|_{\rho_{r+1}=\rho_r} \end{aligned}$$

We again use the Fundamental Theorem of Calculus with respect to  $\rho_{r+1}$  to interpolate between  $\rho_{r+1} = \rho_r$  and  $\rho_{r+1} = 0$ , hence

$$\begin{aligned} \frac{\partial^r f}{\partial s_{l_1} \dots \partial s_{l_r}} (\hat{t}(F, \vec{\rho})) &= \frac{\partial^r f}{\partial s_{l_1} \dots \partial s_{l_r}} (t(F, \vec{\rho})) \\ &+ \int_0^{\rho_r} d\rho_{r+1} \sum_{l_{r+1}} \frac{\partial^{r+1} f}{\partial s_{l_1} \dots \partial s_{l_r} \partial s_{l_{r+1}}} (\hat{t}((F, l_{r+1}), (\vec{\rho}, \rho_{r+1}))) \end{aligned}$$

where the sum is over pairs  $l_{r+1}$  corresponding to the nonzero entries of  $v_{\hat{1}} - v_{\pi_r}$ . This is tantamount to summing over all pairs not falling inside a connected component of  $F$ , i.e., all pairs one can append to  $F$  in order to produce an o-forest of length  $r+1$ . This immediately implies identity (1) for  $r+1$ . Finally, the identity (1) reduces to the statement of Theorem 2 as soon as  $r$  reaches  $|E| - 1$ . This is because when  $r = |E| - 1$  the second sum in (1) is over ordered connecting trees  $F$  for which  $\hat{t}(F, \vec{\rho})$  and  $t(F, \vec{\rho})$  are the same.  $\blacksquare$

**Proof of Theorem 1:** Starting from the identity in Theorem 2 we collect the o-forests corresponding to an unordered forest  $\mathfrak{F} = \{l_1, \dots, l_p\}$ . This contributes

$$\sum_{\sigma \in \mathfrak{S}_p} \int_{1 > \rho_1 > \dots > \rho_p > 0} d\vec{\rho} \frac{\partial^{|\mathfrak{F}|} f}{\prod_{l \in \mathfrak{F}} \partial s_l} (t(F^\sigma, \vec{\rho})) \quad (3)$$

where we ignored some set of measure zero, and used the notation  $F^\sigma = (l_{\sigma(1)}, \dots, l_{\sigma(p)})$ . Now define the family (rather than sequence) of variables  $\vec{h} = (h_l)_{l \in \mathfrak{F}}$  by letting  $h_{\sigma(q)} = \rho_q$  for any  $q$ ,  $1 \leq q \leq p$ . One can check from the previous definitions that

$$t(F^\sigma, \vec{\rho}) = s(\mathfrak{F}, \vec{h})$$

As a result, the quantity (3) becomes

$$\sum_{\sigma \in \mathfrak{S}_p} \int_{1 > h_{l_{\sigma(1)}} > \dots > h_{l_{\sigma(p)}} > 0} d\vec{h} \frac{\partial^{|\mathfrak{F}|} f}{\prod_{l \in \mathfrak{F}} \partial s_l} (s(\mathfrak{F}, \vec{h})) = \int_{[0,1]^{\mathfrak{F}}} d\vec{h} \frac{\partial^{|\mathfrak{F}|} f}{\prod_{l \in \mathfrak{F}} \partial s_l} (s(\mathfrak{F}, \vec{h}))$$

by combining the simplices of integration accounting for the relative ordering of the parameters into the full cube  $[0,1]^{\mathfrak{F}}$ . The statement about the arguments  $s(\mathfrak{F}, \vec{h})$  belonging to  $\mathcal{K}_E$  follows from Prop. 1.  $\blacksquare$

We will now deduce a few lemmas as corollaries of the BKAR forest formula.

**Lemma 1** *Again let us consider a finite set  $E$  and let us denote by  $E^{(2)}$  the set of unordered pairs  $l = \{a, b\}$  in  $E$ . Let  $V_{\{a,b\}}$  be a collection of complex numbers indexed by  $E^{(2)}$ . Then*

$$\sum_{\mathfrak{g} \rightsquigarrow E} \prod_{l \in \mathfrak{g}} (e^{-V_l} - 1) = \sum_{\substack{\mathfrak{T} \rightsquigarrow E \\ \mathfrak{T} \text{ tree}}} \int_{[0,1]^{\mathfrak{T}}} d\vec{h} \left\{ \prod_{l \in \mathfrak{T}} (-V_l) \right\} e^{-\sum_{l \in E^{(2)}} s(\mathfrak{T}, \vec{h})_l V_l}. \quad (4)$$

Here  $\mathbf{g}$  is summed over all simple graphs (i.e. subsets of  $E^{(2)}$ ) which connect  $E$ . We abbreviate this property by the notation  $\mathbf{g} \rightsquigarrow E$ . On the right-hand side the sum is on spanning trees  $\mathfrak{T}$  which connect  $E$ . The notation  $s(\mathfrak{T}, \vec{h})$  is as in Theorem 1.

**Proof:** This is a consequence of Theorem 1 and the uniqueness of the Möbius inverse. Given any nonempty subset  $X \subset E$ , let  $\gamma_1(X)$  denote the expression analogous to the left-hand side of (4) for  $X$  instead of  $E$ . Namely one sums over graphs  $\mathbf{g} \subset X^{(2)}$  connecting  $X$ , and the  $V_l$  are the ones coming from  $E$  by restriction:

$$\gamma_1(X) = \sum_{\mathbf{g} \rightsquigarrow X} \prod_{l \in \mathbf{g}} (e^{-V_l} - 1) .$$

Likewise let  $\gamma_2(X)$  be the expression analogous to the right-hand side of (4) for  $X$  instead of  $E$ . For any partition  $\pi \in \Pi_E$  we define

$$c_1(\pi) = \prod_{X \in \pi} \gamma_1(X)$$

and

$$c_2(\pi) = \prod_{X \in \pi} \gamma_2(X) .$$

Let us also define

$$d(\pi) = \prod_{l \in E^{(2)}} \left[ e^{-\mathbb{1}\{\exists X \in \pi, l \subset X\}} V_l \right] .$$

Let us denote the natural order relation on the partition lattice  $\Pi_E$  by  $\preceq$ , i.e., one writes  $\pi \preceq \pi'$  if partition  $\pi$  is a refinement of partition  $\pi'$ . We will first show that

$$d(\pi) = \sum_{\pi' \preceq \pi} c_1(\pi') .$$

Indeed, writing

$$e^{-\mathbb{1}\{\exists X \in \pi, l \subset X\}} V_l = 1 + \mathbb{1}\{\exists X \in \pi, l \subset X\} (e^{-V_l} - 1)$$

and expanding one has

$$d(\pi) = \sum_{\mathbf{g} \subset E^{(2)}} \mathbb{1} \left\{ \begin{array}{l} \forall l \in \mathbf{g} \\ \exists X \in \pi, l \subset X \end{array} \right\} \prod_{l \in \mathbf{g}} (e^{-V_l} - 1) .$$

Let us denote by  $\pi(\mathbf{g}) \in \Pi_E$  the partition of connected components of a graph  $\mathbf{g} \in E^{(2)}$ . We then have by collecting the outcome of the expansion with respect to the connected components

$$\begin{aligned} d(\pi) &= \sum_{\pi' \in \Pi_E} \sum_{\substack{\mathbf{g} \subset E^{(2)} \\ \pi(\mathbf{g}) = \pi'}} \mathbb{1} \left\{ \begin{array}{l} \forall l \in \mathbf{g} \\ \exists X \in \pi, l \subset X \end{array} \right\} \prod_{l \in \mathbf{g}} (e^{-V_l} - 1) \\ &= \sum_{\pi' \preceq \pi} \prod_{X \in \pi'} \left\{ \sum_{\mathbf{g} \rightsquigarrow X} \prod_{l \in \mathbf{g}} (e^{-V_l} - 1) \right\} \end{aligned}$$

as wanted. Now we also have

$$d(\pi) = \prod_{X \in \pi} \prod_{l \in X^{(2)}} e^{-V_l}.$$

For any such  $X$  we consider the function

$$f_X(s) = \prod_{l \in X^{(2)}} e^{-s_l V_l}$$

for multiplsets  $s = (s_l)_{l \in X^{(2)}}$  and apply Theorem 1 to it, thus obtaining

$$\begin{aligned} \prod_{l \in X^{(2)}} e^{-V_l} &= f_X(1) \\ &= \sum_{\mathfrak{F}_X \text{ forest on } X} \int_{[0,1]^{\mathfrak{F}_X}} d\vec{h}_X \frac{\partial^{|\mathfrak{F}_X|} f_X}{\prod_{l \in \mathfrak{F}_X} \partial s_l} \left( s(\mathfrak{F}_X, \vec{h}_X) \right) \\ &= \sum_{\mathfrak{F}_X \text{ forest on } X} \int_{[0,1]^{\mathfrak{F}_X}} d\vec{h}_X \left\{ \prod_{l \in \mathfrak{F}_X} (-V_l) \right\} e^{-\sum_{l \in X^{(2)}} s(\mathfrak{F}_X, \vec{h}_X)_l V_l}. \end{aligned}$$

Now again collecting the terms componentwise, one can rewrite the latter expression as

$$\prod_{l \in X^{(2)}} e^{-V_l} = \sum_{\pi_X \in \Pi_X} \prod_{Y \in \pi_X} \left[ \sum_{\substack{\mathfrak{F}_Y \rightsquigarrow Y \\ \mathfrak{F}_Y \text{ tree}}} \int_{[0,1]^{\mathfrak{F}_Y}} d\vec{h}_Y \left\{ \prod_{l \in \mathfrak{F}_Y} (-V_l) \right\} e^{-\sum_{l \in Y^{(2)}} s(\mathfrak{F}_Y, \vec{h}_Y)_l V_l} \right].$$

This is because of the definition of the  $s(\mathfrak{F}, \vec{h})_l$  in the BKAR formula. These vanish for pairs which are not inside a connected component. Whereas for pairs  $l$  which are inside a connected component, the  $s(\mathfrak{F}, \vec{h})_l$  only depend on the edges of the forest  $\mathfrak{F}$  which are in that component. Now

$$\prod_{l \in X^{(2)}} e^{-V_l} = \sum_{\pi_X \in \Pi_X} \prod_{Y \in \pi_X} \gamma_2(Y)$$

and as a result

$$\begin{aligned} d(\pi) &= \prod_{X \in \pi} \left( \sum_{\pi_X \in \Pi_X} \prod_{Y \in \pi_X} \gamma_2(Y) \right) \\ &= \sum_{\pi' \preceq \pi} \prod_{Y \in \pi'} \gamma_2(Y) \end{aligned}$$

where we collected the blocks of the partitions  $\pi_X$ , for  $X$  a block of  $\pi$ , into a single partion  $\pi'$  of  $E$  which is refinement of  $\pi$ . Hence

$$d(\pi) = \sum_{\pi' \preceq \pi} c_1(\pi')$$

Therefore both  $c_1$  and  $c_2$  are Möbius inverses of  $d$  on the partition lattice  $\Pi_E$ , and they must be equal. Specializing to  $c_1(\hat{1}) = c_2(\hat{1})$  proves the lemma.  $\blacksquare$

The following tree graph inequality, initially due to Brydges, Battle and Federbush (see [5, 4, 18]) is useful when performing Mayer expansions for a gas of particles with unbounded interaction potential energies.

**Lemma 2** *Under the same hypotheses as in Lemma 1, let us assume that the numbers  $V_l$  satisfy, in addition, the following stability hypothesis: there are nonnegative numbers  $U_a$ , for  $a \in E$ , such that for any subset  $S \subset E$  one has*

$$\left| \sum_{l \in S^{(2)}} V_l \right| \leq \sum_{a \in S} U_a .$$

Then the following inequality holds

$$\left| \sum_{\mathfrak{g} \rightsquigarrow E} \prod_{l \in \mathfrak{g}} (e^{-V_l} - 1) \right| \leq e^{\sum_{a \in E} U_a} \sum_{\substack{\mathfrak{T} \rightsquigarrow E \\ \mathfrak{T} \text{ tree}}} \prod_{l \in \mathfrak{T}} |V_l| .$$

**Proof:** This is an easy consequence of Lemma 1 and the fact that  $s(\mathfrak{T}, \vec{h})$  appearing in the BKAR formula is in  $\mathcal{K}_E$ , i.e., is a convex combination of partition vectors. Indeed, for any given  $\mathfrak{T}$  and  $\vec{h}$  as in (4), one can find nonnegative numbers  $\lambda_1, \dots, \lambda_p$ , satisfying  $\sum_{q=1}^p \lambda_q = 1$ , as well as partitions  $\pi_1, \dots, \pi_p \in \Pi_E$  such that

$$s(\mathfrak{T}, \vec{h}) = \sum_{q=1}^p \lambda_q v_{\pi_q} .$$

Therefore

$$\begin{aligned} e^{-\sum_{l \in E^{(2)}} s(\mathfrak{T}, \vec{h})_l V_l} &= \exp \left[ - \sum_{q=1}^p \lambda_q \sum_{l \in E^{(2)}} (v_{\pi_q})_l V_l \right] \\ &= \exp \left[ - \sum_{q=1}^p \lambda_q \sum_{X \in \pi_q} \sum_{l \in X^{(2)}} V_l \right] . \end{aligned}$$

From which one derives

$$\begin{aligned} \left| e^{-\sum_{l \in E^{(2)}} s(\mathfrak{T}, \vec{h})_l V_l} \right| &\leq \exp \left[ \sum_{q=1}^p \lambda_q \sum_{X \in \pi_q} \left| \sum_{l \in X^{(2)}} V_l \right| \right] \\ &\leq \exp \left[ \sum_{q=1}^p \lambda_q \sum_{X \in \pi_q} \sum_{a \in X} U_a \right] \\ &\leq \exp \left[ \sum_{a \in E} U_a \right] \end{aligned}$$

using the stability hypothesis. Now the desired inequality clearly follows from the formula (4).  $\blacksquare$

**Lemma 3** *Under the same hypotheses as in lemma 1 we have*

$$\sum_{\mathbf{g} \rightsquigarrow E} \prod_{l \in \mathbf{g}} (-V_l) = \sum_{\substack{\mathfrak{T} \rightsquigarrow E \\ \mathfrak{T} \text{ tree}}} \int_{[0,1]^{\mathfrak{T}}} d\vec{h} \prod_{l \in \mathfrak{T}} (-V_l) \prod_{l \in E^{(2)} \setminus \mathfrak{T}} (1 - s(\mathfrak{T}, \vec{h})_l V_l) .$$

**Proof:** The proof follows the same lines as that of Lemma 1. This time

$$d(\pi) = \prod_{l \in E^{(2)}} [\mathbb{1}\{\exists X \in \pi, l \subset X\} (1 - V_l)]$$

and the function to which one applies the BKAR formula is

$$f_X(s) = \prod_{l \in X^{(2)}} (1 - s_l V_l) .$$

The rest of the argument based on the uniqueness of the Möbius inverse is the same. ■

An immediate corollary is the so-called Penrose-Rota inequality which bounds Mayer coefficients by a sum over spanning trees [14, 20].

**Some history:** The formula discovered by Brydges and Kennedy [8] is the one stated in Lemma 1. They used 1 minus the parameters, so the coefficients of the  $V$ 's in the exponential involve a maximum over the connecting path instead of a minimum. Their proof uses an explicit tree-sum solution for a differential equation of Hamilton-Jacobi type (see [6] for a nice presentation). It is inspired by the Wilson-Polshinski continuous renormalization group differential equation [22, 16]. The reason why such solutions of nonlinear differential equations can be expressed as sums over trees is explained from a combinatorial point of view, e.g., in [13]. The fundamental calculus version of Theorem 1 first appeared in [1]. It was there called the Brydges-Kennedy Taylor forest formula. D. Brydges calls it the Abdesselam-Rivasseau formula in his lecture at the conference “Combinatorial Identities and Their Applications in Statistical Mechanics”, Cambridge, April 2008:

<http://www.newton.ac.uk/programmes/CSM/csmw03.html>

In [19] the name Brydges-Kennedy-Abdesselam-Rivasseau formula was first used. The proof given in [1] is purely algebraic and relies on a combinatorial partial fraction decomposition identity. Two other proofs of this partial fraction identity by A. Abdesselam and V. Lafforgue can be found at:

<http://people.virginia.edu/aa4cr/forestpage.html>

The proof given here originates from ideas of H. Knörrer, J. Magnen and V. Rivasseau. It was considerably generalized in [2] to expansions which allow hypergraphs and p-PI connectivity. A similar generalization was independently discovered by G. Poirot [15].

## References

- [1] A. Abdesselam and V. Rivasseau. Trees, forests and jungles: a botanical garden for cluster expansions. Constructive physics (Palaiseau, 1994), 7–36, Lecture Notes in Phys., 446. Berlin: Springer, 1995.



- [2] A. Abdesselam and V. Rivasseau. An explicit large versus small field multiscale cluster expansion. *Rev. Math. Phys.* 9 (1997), no. 2, 123–199.
- [3] R. A. Adams and J. J. F. Fournier. *Sobolev Spaces*. Second edition. New York: Academic Press, 2003.
- [4] G. A. Battle and P. Federbush. A note on cluster expansions, tree graph identities, extra  $1/N!$  factors! *Lett. Math. Phys.* 8 (1984), no. 1, 55–57.
- [5] D. C. Brydges. A short course on cluster expansions. *Phénomènes critiques, systèmes aléatoires, théories de jauge, Part I, II* (Les Houches, 1984), 129–183, North-Holland, Amsterdam, 1986.
- [6] D. C. Brydges. *Functional Integrals and their Applications* (Notes for a course for the “Troisième Cycle de la Physique en Suisse Romande” given in Lausanne, Switzerland, during the summer of 1992). Notes with the collaboration of R. Fernández.  
  
[http://www.ma.utexas.edu/mp\\_arc-bin/mpa?yn=93-24](http://www.ma.utexas.edu/mp_arc-bin/mpa?yn=93-24)
- [7] D. C. Brydges and P. Federbush. A new form of the Mayer expansion in classical statistical mechanics. *J. Math. Phys.* 19 (1978), no. 10, 2064–2067.
- [8] D. Brydges and T. Kennedy. Mayer expansions and the Hamilton-Jacobi equation. *J. Stat. Phys.* 48 (1987), 19
- [9] D. Brydges and P. Martin. Coulomb systems at low density: a review. *J. Stat. Phys.* 96 (1999), 1163–1330.
- [10] J. Glimm and A. Jaffe. *Quantum physics. A functional integral point of view*. Second edition. New York: Springer, 1987.
- [11] J. Glimm, A. Jaffe and T. Spencer. The Wightman axioms and particle structure in the  $P(\phi)_2$  quantum field model. *Annals of Math.* 100 (1974), 585–632.
- [12] J. Glimm, A. Jaffe and T. Spencer. The particle structure of the weakly coupled  $P(\phi)_2$  model and other applications of high temperature expansions, Part II: The cluster expansion. *Constructive quantum field theory, Erice 1973*, eds. G. Velo and A. Wightman, *Lecture notes in Phys.* 25. New York: Springer 1973.
- [13] P. Leroux and G. X. Viennot. Combinatorial resolution of systems of differential equations. I. Ordinary differential equations. *Combinatoire énumérative* (Montreal, Que., 1985/Quebec, Que., 1985), 210–245, *Lecture Notes in Math.*, 1234, Springer, Berlin, 1986.
- [14] O. Penrose. Convergence of fugacity expansions for classical systems. In “*Statistical Mechanics, Foundations and Applications*”, T. A. Bak, ed., IUPAP Meeting, Copenhagen, 1966. Benjamin Inc., New York, 1967, pp. 101–109.

- [15] G. Poirot. Mean Green's function of the Anderson model at weak disorder with an infra-red cut-off. *Ann. Inst. H. Poincaré Phys. Théor.* 70 (1999), no. 1, 101–146.
- [16] J. Polchinski. Renormalization and effective lagrangians. *Nucl. Phys. B* 231 (1984), 269–295.
- [17] A. Pordt. Mayer expansions for Euclidean lattice field theory: convergence properties and relation with perturbation theory. Desy preprint 85-103, unpublished, 1985. Available at <http://www-lib.kek.jp/top-e.html>
- [18] A. Procacci, B. N. B. de Lima and B. Scoppola. A remark on high temperature polymer expansion for lattice systems with infinite range pair interactions. *Lett. Math. Phys.* 45 (1998), no. 4, 303–322.
- [19] V. Rivasseau. Lectures at the summer school “Rigorous Results in Statistical Mechanics and Quantum Field Theory”, June 2008, Feza Gürsey Institute, Istanbul. <http://www.gursey.gov.tr/semesters/mathphys08/notes/rivasseau-lec2.pdf>
- [20] G.-C. Rota. On the foundations of combinatorial theory. I. Theory of Möbius functions. *Z. Wahrscheinlichkeitstheorie und Verw. Gebiete* 2 1964 340–368 (1964).
- [21] W. Wagner. Analyticity and Borel-summability of the perturbation expansion for correlation functions of continuous spin systems. *Helv. Phys. Acta* 54 (1981/82), no. 3, 341–363.
- [22] K. G. Wilson and J. Kogut. Renormalisation group and the  $\epsilon$  expansion. *Phys. Rep.* 12 (1974), 75–200.